

Mechanism of Exact Transition between Cationic and Anionic Redox Activities in Cathode Material $\text{Li}_2\text{FeSiO}_4$

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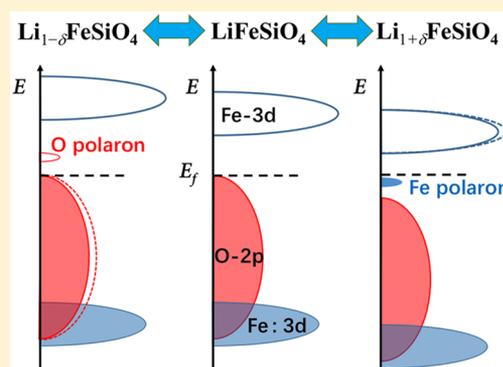
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S Supporting Information

ABSTRACT: The discovery of anion redox activity is promising for boosting the capacity of lithium ion battery (LIB) cathodes. However, fundamental understanding of the mechanisms that trigger the anionic redox is still lacking. Here, using hybrid density functional study combined with experimental soft X-ray absorption spectroscopy (sXAS) measurements, we unambiguously proved that $\text{Li}_{(2-x)}\text{FeSiO}_4$ performs sequent cationic and anionic redox activity through delithiation. Specifically, Fe^{2+} is oxidized to Fe^{3+} during the first Li ion extraction per formula unit (f.u.), while the second Li ion extraction triggered the oxygen redox exclusively. Cationic and anionic redox result in electron and hole polaron states, respectively, explaining the poor conductivity of $\text{Li}_{(2-x)}\text{FeSiO}_4$ noted by previous experiments. In contrast, other cathode materials in this family exhibit diversity of the redox process. $\text{Li}_2\text{MnSiO}_4$ shows double cationic redox ($\text{Mn}^{2+}-\text{Mn}^{4+}$) during the whole delithiation, while $\text{Li}_2\text{CoSiO}_4$ shows simultaneous cationic and anionic redox. The present finding not only provides new insights into the oxygen redox activity in polyanionic compounds for rechargeable batteries but also sheds light on the future design of high-capacity rechargeable batteries.



Charge compensation of transition metals (TMs) in TM-based compounds when the total charge of the compound is altered (e.g., through oxidation/reduction, carrier injection, chemical doping) is a fundamental and intriguing topic in both physics and chemistry.^{1,2} Classic inorganic chemistry tacitly assumes that most of the changes of the total charge are accommodated by a change in the charge of the TM ion.^{3,4} Typically, for rechargeable lithium ion battery (LIB) cathode materials, TM ions were regarded as the sole sources of electrochemical activity in an intercalation cathode to provide charge-compensating electrons after Li ion extraction.^{5–9} As a result, the theoretical capacity of these oxides has been limited by the number of electrons offered by the TM redox reaction and its relatively high atomic weight.

Until recently, this scenario was challenged by the discovery of anionic redox activity in Li-rich layered TM oxides.^{10–19} Additionally, the discovery of anionic redox activity has been extended to other kinds of cathode materials beyond Li-rich layered TM oxides, such as nonlayered TM oxides^{20,21} and metal–organic compounds.²² Anionic reduction not only extends our understanding of the charge compensation process during delithiation in cathodes but also raises the opportunity to boost the capacity and energy density of LIBs by combining

both cationic and anionic redox processes within the same material.^{14,23,24}

However, such staggering capacities suffer from capacity fade with cycling,²⁵ which is mainly due to the irreversible loss of lattice oxygen during the anionic redox process.^{15,19,26} Thus, clarifying the specific redox mechanism and fundamental understanding of the local structure and electronic state evolutions during the anionic redox process become crucial for both experimentalists and theorists. In general, three types of redox processes are revealed in the reported cathode materials with anionic redox, including (1) cationic redox first and anionic redox next, such as in layered NMC,^{15,27,28} (2) simultaneous cationic and anionic redox, such as in Li_5FeO_4 ,²⁰ and (3) anionic redox only, such as in Li_2MnO_3 .^{29–31} Here, we focus on a series of silicate cathode materials $\text{Li}_{(2-x)}\text{TMSiO}_4$ (TM = Mn, Fe, Co, and Ni), for which the mechanism of redox activity is still under debate.^{32–39} Using hybrid density functional study combined with the modern perspective of doping and experimental soft X-ray absorption spectroscopy

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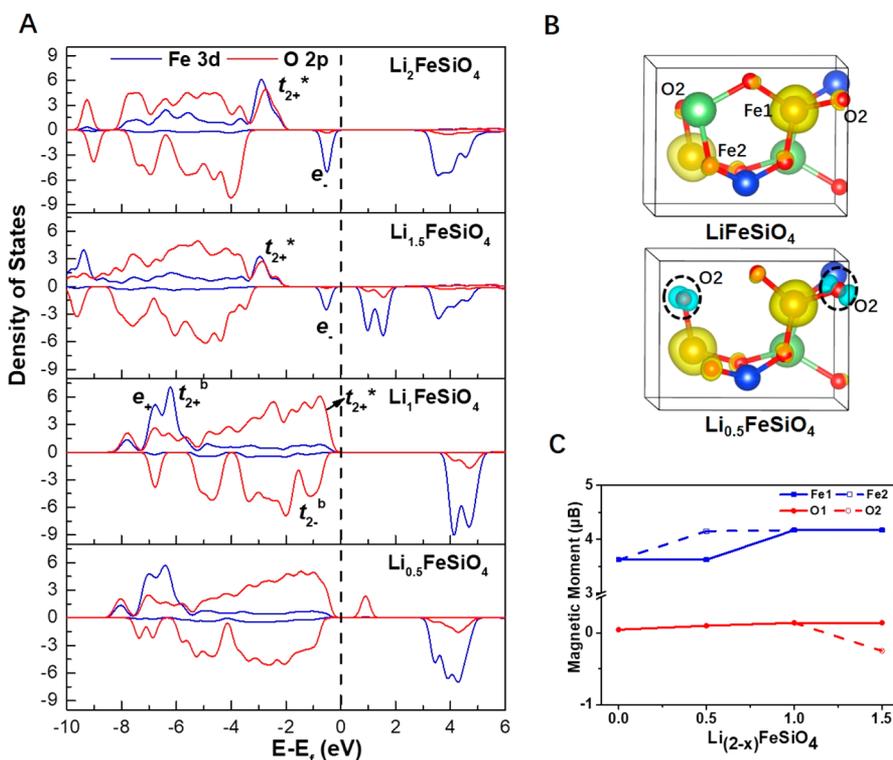


Figure 1. HSE calculated electronic structures for $\text{Li}_{(2-x)}\text{FeSiO}_4$ during delithiation. (A) Projected density of states (PDOS) of Fe-3d and O-2p states in $\text{Li}_{(2-x)}\text{FeSiO}_4$ ($x = 0, 0.5, 1.0$, and 1.5). (B) Isosurface of the spin density in the unit cell of partially delithiated $\text{Li}_{(2-x)}\text{FeSiO}_4$ ($x = 1.0$ and 1.5). Green, Li; red, O; dark blue, Si; brown, Fe. The isovalues for the isosurfaces of the spin densities were $0.05 \text{ e} \text{ \AA}^{-3}$ for all cases. (C) Calculated magnetic moment of Fe and O in $\text{Li}_{(2-x)}\text{FeSiO}_4$ at different delithiation states. O1 and O2 denote two kinds of O ions with different variations of the magnetic moment in $\text{Li}_{(2-x)}\text{FeSiO}_4$. O2 are denoted in (B), and all of the remaining unidentified O ions in (B) are O1.

(sXAS) measurements, we proved that $\text{Li}_{(2-x)}\text{FeSiO}_4$ ($0 < x < 2$) performs sequent cation and anionic redox activity upon delithiation. During delithiation, our results demonstrated that Fe^{2+} ions change to Fe^{3+} until a critical Li concentration $x = 1$, while further Li extraction triggers oxygen redox instead of cation oxidation. In order to uncover the physical origin of the exact transition point ($x = 1$) of cationic/anionic redox, we then focused on dilute lithiation and delithiation behaviors of LiFeSiO_4 and found that both of lithiation and delithiation processes result in a polaronic state, but with cationic and anionic character, respectively. In contrast to the previous reported oxygen clustering during oxygen redox in Li-rich layered TM oxides, the oxygen polaron in $\text{Li}_{(2-x)}\text{FeSiO}_4$ is localized on a single O ion, which avoids peroxy-like O–O bond formation to prevent the lattice O loss. Moreover, we found that not all of the polyoxoanion $\text{Li}_{(2-x)}\text{TMSiO}_4$ exhibit such polaronic oxygen redox activity after more than one Li ion is extracted but show a strong dependence on the 3d occupation numbers of TM elements.

In the polyoxoanion-type intercalation $\text{Li}_{(2-x)}\text{FeSiO}_4$ structures, each iron, lithium, and silicon atom is coordinated by four oxygen atoms, while the FeO_4 tetrahedra are cross-linked by silicate groups (Figure S1). The electrochemical performance of $\text{Li}_2\text{FeSiO}_4$ is discussed in Supporting Information Note 1. The first discharge specific capacity is 305 mAh/g , corresponding to 1.86 Li-storages per $\text{Li}_2\text{FeSiO}_4$ unit with oxygen redox occurring, close to the previously reported performance of $\text{Li}_2\text{FeSiO}_4$ nanocrystals.³⁴ To reveal the underlying mechanism of the redox process during delithiation, we perform DFT calculations to revisit different delithiation stages of $\text{Li}_{(2-x)}\text{FeSiO}_4$ and compare the results with our sXAS

measurements. The general challenge of standard exchange–correlation functionals in DFT is the notorious self-interaction error (SIE). Therefore, previous DFT studies even predicted a metallic state for $\text{Li}_{(2-x)}\text{FeSiO}_4$ during delithiation,^{32,40} in sharp contrast to the semiconducting phase as reported by experiments.^{41,42}

Sequent Cation and Anion Redox in $\text{Li}_{(2-x)}\text{FeSiO}_4$. We first apply a hybrid functional in the form of Heyd–Scuseria–Ernzerhof (HSE06)⁴³ to correct the SIE. The evolution of the density of states for the unit cell ($\text{Li}_4\text{Fe}_2\text{Si}_2\text{O}_8$) of partially delithiated $\text{Li}_{(2-x)}\text{FeSiO}_4$ ($x = 0, 0.5, 1.0$, and 1.5) is shown in Figure 1A (the corresponding crystal structures is shown in Figure S3). It is found that when x goes from 0 to 1.5, the band gap E_g changes from 2.91 ($x = 0$) to 3.60 and 0.57 eV for $x = 1$ and 1.5 , respectively, consistent with the semiconducting phase as reported by experiments.^{41,42} Moreover, our results naturally explain that $\text{Li}_{(2-x)}\text{FeSiO}_4$ always keeps a poor electronic conductivity and rate capability during electrochemical cycling.⁴⁴ From the top three panels of Figure 1A, we can see that the Fe-3d peak right below the Fermi level is gradually shifted upward after the first Li extraction per f.u., indicating the redox couple of $\text{Fe}^{2+}/\text{Fe}^{3+}$. However, the second Li ion extraction shifts valence states with pure O-2p character across the Fermi level (Figure 1A, bottom panel), indicating that instead of $\text{Fe}^{3+}/\text{Fe}^{4+}$ it is an O anion playing an active role. Interestingly, the upward shifted O-2p states form a localized, polaron-like hole state inside of the band gap of LiFeSiO_4 . Such a hole-localized state is confirmed by the isosurface of the spin density of partially delithiated $\text{Li}_{(2-x)}\text{FeSiO}_4$ ($x = 1$ and 1.5), as shown in Figure 1B. We observe large spin density from two oxygen ions with the shape of an isolated O-2p

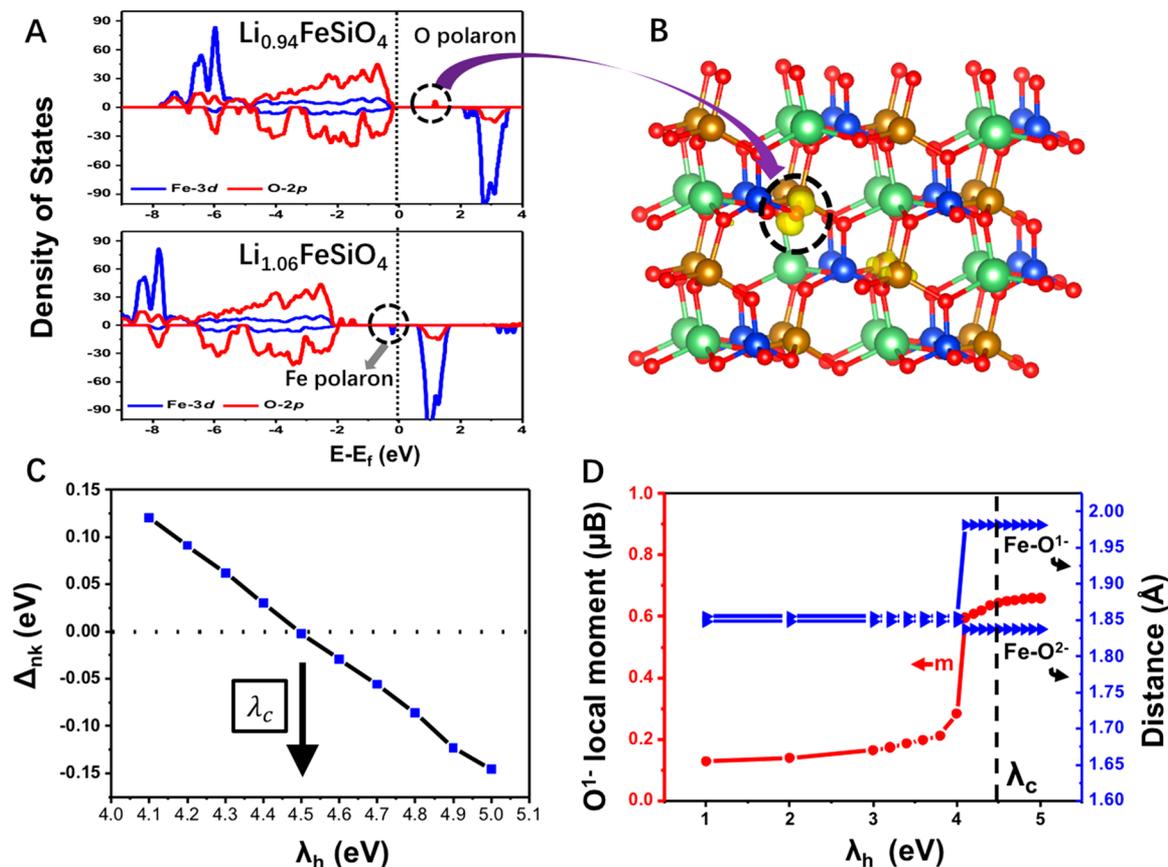


Figure 2. Cation/anion transition at the critical point $x = 1$. (A) PDOS of Fe-3d and O-2p states in $\text{Li}_{0.94}\text{FeSiO}_4$ (delithiation, top panel) and $\text{Li}_{1.06}\text{FeSiO}_4$ (lithiation, bottom panel). (B) Isosurface of the charge density of the O polaron state in the cell of $\text{Li}_{15}\text{Fe}_{16}\text{Si}_{16}\text{O}_{64}$, showing localization on a single O ion. (C) Non-Koopmans energy Δ_{nk} (see eq 1) as a function of the hole-state potential strength λ_h (see eq 2). (D) Structural and magnetic properties of the O^- in $\text{Li}_{15}\text{Fe}_{16}\text{Si}_{16}\text{O}_{64}$ as a function of the hole-state potential strength λ_h . The polaronic state is stable above a value $\lambda_h > 4.0$ eV, while the critical value λ_c that fulfills the generalized Koopmans theory is marked by the dashed line. m : local magnetic moment of O^- .

orbital for $\text{Li}_{0.5}\text{FeSiO}_4$, indicating that this polaron state is localized on two O ions belonging to one SiO_4 tetrahedron. The calculated magnetic moments of Fe and O in partially delithiated $\text{Li}_{(2-x)}\text{FeSiO}_4$ (Figure 1C) also support the whole redox process analyzed above. From $\text{Li}_2\text{FeSiO}_4$ to LiFeSiO_4 , the magnetic moments of the two Fe cations change from $3.61 \mu\text{B}$ (Fe^{2+} , $e_g^3 t_{2g}^3$) to $4.17 \mu\text{B}$ (Fe^{3+} , $e_g^2 t_{2g}^3$) by sequence. For the further delithiation (from $x = 1$ to 1.5), the magnetic moments of Fe ions remain nearly unchanged, while the magnetic moments of two O ions (unchanged from $\text{Li}_2\text{FeSiO}_4$ to LiFeSiO_4) start to change to $-0.25 \mu\text{B}$ in $\text{Li}_{0.5}\text{FeSiO}_4$. The structural response to the creation of hole polarons is discussed in Supporting Information Note 2. We find that two adjacent O ions form localization centers of the isolated state for $\text{Li}_{0.5}\text{FeSiO}_4$, indicating a $2\text{O}^{2-}/(\text{O}_2)^{3-}$ process, which is similar to the peroxo-like oxygen redox process in layered TM oxides.^{11,16–19,25} To further confirm the above theoretical calculations, the Fe L-edge and O K-edge sXAS of samples $\text{Li}_2\text{FeSiO}_4$, LiFeSiO_4 , and $\text{Li}_{0.15}\text{FeSiO}_4$ are tested, as shown in Figure S5A,B, respectively. Our sXAS results provide consistent evidence for the physical picture suggested by our theoretical calculations. (see Supporting Information Note 3 for details).

Cation/Anion Transition at the Critical Point $x = 1$. A general discussion of the origin of anionic redox activity in $\text{Li}_{(2-x)}\text{FeSiO}_4$ can be found in Supporting Information Note

4. In the following, to get a microscopic picture of how the switch between cation and anion redox is triggered at the transition point $x = 1$ for $\text{Li}_{(2-x)}\text{FeSiO}_4$, we focused on lithiation and delithiation behaviors of LiFeSiO_4 at a small concentration, which corresponds to a doping perspective. The challenge of such prediction of dilute doping lies in that especially for delithiation the localized hole states are often incorrectly described in standard DFT exchange–correlation functionals as rather delocalized states that spread over all oxygen ligands.^{45–47} Such “delocalization error”⁴⁸ originates from the convex bowing of the total energy $E(N)$ with respect to electron occupation number N . Indeed, a perfect correction is required to fulfill the so-called generalized Koopmans condition

$$\Delta_{nk} = E(N + 1) - E(N) + \text{eig}(N) = 0 \quad (1)$$

where $E(N + 1) - E(N)$ denotes the total energy cost to add an electron from the hole-doped system and $\text{eig}(N)$ is the single-particle energy of the lowest unoccupied state in the electron-doped system. Here we use the Lany–Zunger approach⁴⁹ to restore the generalized Koopmans condition (see the Supporting Information for details). Basically, we introduce an onsite potential implemented as the combination of DFT+U and nonlocal external potentials⁵⁰ to certain atomic orbitals. For hole doping, the corresponding electron state potential is given by

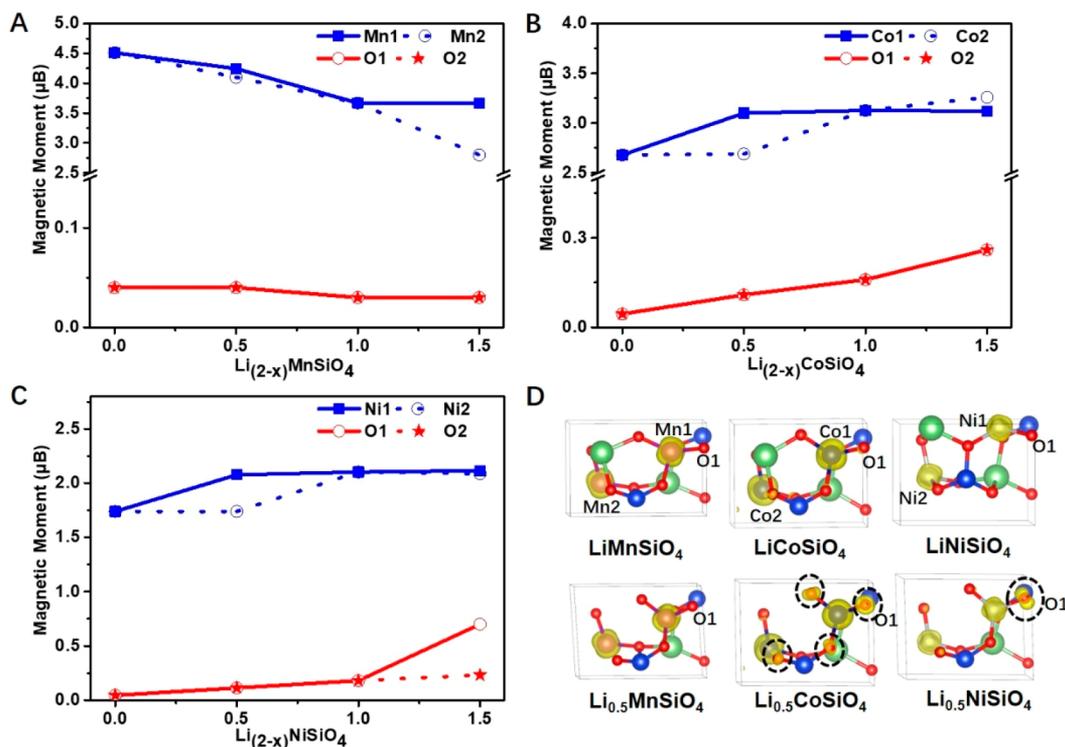


Figure 3. Redox activity in other $\text{Li}_{(2-x)}\text{TMSiO}_4$ compounds. (A–C) Calculated magnetic moment of TM and O in $\text{Li}_{(2-x)}\text{TMSiO}_4$ (TM = Mn, Co, and Ni) at different delithiation states. O1 are denoted in (D), and all of the remaining unidentified O ions in (D) are O2. (D) Isosurface of the spin density in the unit cell of partially delithiated $\text{Li}_{(2-x)}\text{TMSiO}_4$ ($x = 1.0$ and 0.5). Green, Li; red, O; dark blue, Si; purple, Mn; light blue, Co; silver, Ni.

$$V_h = \lambda_h(1 - n_{m,\sigma}/n_{\text{host}}) \quad (2)$$

where $n_{m,\sigma}$ and n_{host} denote the occupation of the m sublevel of spin σ and the occupation of the host material without doping, respectively. We determine n_{host} in eq 2 to be 0.595 from the partial charge of the O-p orbitals of the occupied states. The parameter λ_h is tuned to perfectly fulfill the linearity of $E(N)$, i.e., the generalized Koopmans condition eq 1.

Starting from the cation/anion transition point ($x = 1$), we use a $2 \times 2 \times 2$ supercell (16 Li atoms per cell) of LiFeSiO_4 and remove or add one Li atom to achieve a relatively low carrier doping concentration ($\sim 6\%$). By removing one Li atom, the projected density of states (PDOS) of $\text{Li}_{0.94}\text{FeSiO}_4$ (Figure 2A, top panel) exhibits a localized O acceptor state inside of the fundamental band gap between the Fe-like Hubbard band and O-like charge-transfer band, similar to $\text{Li}_{0.5}\text{FeSiO}_4$ (Figure 1A). Note that here we choose λ_h to be the critical value $\lambda_c = 4.5$ eV, at which point the generalized Koopmans condition eq 1 is exactly fulfilled, i.e., $\Delta_{\text{nk}} = 0$ (Figure 2C). However, the charge density of the localized state (Figure 2B) shows that the small oxygen polaron in $\text{Li}_{0.94}\text{FeSiO}_4$ is localized on a single O ion, rather than an $(\text{O}_2)^{3-}$ cluster for high delithiation concentrations. Such a scenario of a single-O polaron is further verified by the structural and magnetic properties around the Li vacancy. As shown in Figure 2D, when the hole-state potential strength exceeds a critical value of 4.0 eV, a breaking of the tetrahedral symmetry occurs, evidenced by the fact that one Fe–O[−] bond is apparently elongated compared with the other three Fe–O^{2−} bonds and so are the Si–O bonds. In contrast, a standard DFT+U calculation ($\lambda_h = 0$) or small λ_h predicts a high-symmetry extended acceptor state with an equal amplitude of all Fe–O

and Si–O bonds. The polaronic state also suggests the emergence of a local magnetic moment at the O[−] ion that traps the hole. Because our λ_h that fulfills eq 1 (4.5 eV) is higher than the transition value (4.0 eV), the stabilized single-O polaronic state is predicted to be the physical reality in the delithiation process with low concentration. This means that in LiFeSiO_4 delithiation starts by forming a single-O[−] polaron, while the O clusters begin to form when an amount of elongated Fe–O bonds is close enough at a relatively high concentration (e.g., $x = 1.5$).

On the other hand, adding one Li to LiFeSiO_4 (leading to $\text{Li}_{1.06}\text{FeSiO}_4$) exhibits an electron polaron state dominated by Fe-3d orbitals, as shown in Figure 2A (bottom panel). As a result, associated with this lithiation process, an Fe³⁺ ion is reduced to Fe²⁺, corresponding to traditional cation redox activity. Note that we did not apply any onsite potential on the donor state here because, similar to other Fe³⁺ system such as Fe₂O₃, DFT+U already satisfies the generalized Koopmans condition in good approximation.⁵¹ Overall, by a dilute doping approach, we observe exact cation/anion redox transition at LiFeSiO_4 ($x = 1$), with electron/hole single-polaron formation for the cation/anion redox process.

Redox Activity in Other $\text{Li}_{(2-x)}\text{TMSiO}_4$ Compounds. We next consider other $\text{Li}_{(2-x)}\text{TMSiO}_4$ (TM = Mn, Co, and Ni) materials^{32,52–56} to see whether the existence of oxygen redox activity is universal among these polyanionic materials. For $\text{Li}_{(2-x)}\text{MnSiO}_4$, PDOS plots show that upon delithiation from $x = 0$ up to $x = 1.5$ the cation valence $\text{Mn}^{2+} \rightarrow \text{Mn}^{3+} \rightarrow \text{Mn}^{4+}$ accounts for the redox activity (Figure S8). It should be noted that there is also a small portion of oxygen redox activity accompanied by the main Mn redox. The calculated magnetic moments of Mn in partially delithiated $\text{Li}_{(2-x)}\text{MnSiO}_4$ also

support the redox of $\text{Mn}^{2+} \rightarrow \text{Mn}^{3+} \rightarrow \text{Mn}^{4+}$, with the magnetic moment of Mn changing from 4.51 μB in $\text{Li}_2\text{MnSiO}_4$ to 3.69 μB in LiMnSiO_4 and finally to 2.83 μB in $\text{Li}_{0.5}\text{MnSiO}_4$ (Figure 3A). In contrast, the magnetic moment of O remains nearly unchanged, indicating that the small portion of oxygen redox activity observed in the PDOS can be neglected (Figure 3A). Similar to $\text{Li}_2\text{MnSiO}_4$, $\text{Li}_2\text{CoSiO}_4$ also exhibits dominant cationic redox accompanied by a little degree of oxygen redox activity during the whole delithiation process (Figure S9). The only difference is that the oxygen redox in $\text{Li}_2\text{CoSiO}_4$ cannot be neglected, as indicated from Figure 3B that the magnetic moments of both O1 and O2 in $\text{Li}_2\text{CoSiO}_4$ increase from 0.04 to 0.26 μB from $\text{Li}_2\text{CoSiO}_4$ to $\text{Li}_{0.5}\text{CoSiO}_4$. While going to $\text{Li}_2\text{NiSiO}_4$, the redox process resembles the case of $\text{Li}_2\text{FeSiO}_4$ again: a $\text{Ni}^{2+}/\text{Ni}^{3+}$ redox couple for the first Li ion extraction (per f.u.) and a polaron-like oxygen redox couple for the second Li ion extraction (Figure S10). This is also supported by the calculated magnetic moments of Ni in partially delithiated $\text{Li}_{(2-x)}\text{NiSiO}_4$ (Figure 3C), changing from 1.73 μB (Ni^{2+} , $e_g^{4t_{2g}}$) in $\text{Li}_2\text{NiSiO}_4$ to 2.11 μB (Ni^{3+} , $e_g^{4t_{2g}^3}$) in LiNiSiO_4 and then remaining nearly unchanged (2.11 μB for Ni in $\text{Li}_{0.5}\text{NiSiO}_4$). On the other hand, the magnetic moments of one O ion (O1) increase a little by 0.13 μB from $\text{Li}_2\text{NiSiO}_4$ to LiNiSiO_4 and then sharply change to -0.69 μB in $\text{Li}_{0.5}\text{NiSiO}_4$ (Figure 3C).

The isosurfaces of the spin density for the partially delithiated $\text{Li}_{(2-x)}\text{TMSiO}_4$ ($x = 1$ and 1.5) further confirm the above calculations: We observe no spin density for all of the oxygen ions in $\text{Li}_{(2-x)}\text{MnSiO}_4$, nonzero spin density for all of the oxygen ions in $\text{Li}_{(2-x)}\text{CoSiO}_4$, and large spin density with the shape of an isolated O-2p orbital on O1 in $\text{Li}_{0.5}\text{NiSiO}_4$ (Figure 3D). From the above results, we can conclude that the redox activity in $\text{Li}_{(2-x)}\text{TMSiO}_4$ (TM = Mn, Fe, Co, and Ni) shows a strong dependence on the 3d occupation numbers of the TM. This finding recalls the recent observation where a similar oxygen-polaronic effect was found in another type of polyanionic material $\text{Li}_3\text{Fe}_2(\text{PO}_4)_3$ during delithiation, while in $\text{Li}_3\text{V}_2(\text{PO}_4)_3$, the oxidation takes place on the vanadium ions.⁵⁷ A detailed discussion on the evolution of energy orbitals during delithiation explaining the above differences of the redox activity for the $\text{Li}_{(2-x)}\text{TMSiO}_4$ family is provided in Supporting Information Notes 5 and 6.

To summarize, by systematical study of the delithiation process of $\text{Li}_{(2-x)}\text{FeSiO}_4$ with the combination of first-principle calculations and sXAS measurements, we demonstrated a reversible sequent cationic and anionic redox process in $\text{Li}_{(2-x)}\text{FeSiO}_4$. Namely, the first Li ion extraction (per f.u.) is attributed to the cation valence change from Fe^{2+} to Fe^{3+} , while the second Li ion extraction activates almost exclusively the oxygen redox. This full delithiation is reversible and stable until nearly 1.8 Li ions are extracted, corresponding to an extra capacity of 0.8 Li ions induced by anionic redox activity. We then focused on slight lithiation and delithiation behaviors of LiFeSiO_4 from a doping perspective and demonstrated the exact transition point of cation/anion redox at LiFeSiO_4 ($x = 1$). Interestingly, both the lithiation and delithiation processes around the transition point result in electron and hole polaronic states with cationic and anionic character, respectively. In contrast to the previously reported oxygen clustering in Li-rich layered TM oxides, the oxygen polaron in $\text{Li}_{(2-x)}\text{FeSiO}_4$ is localized on a single O ion when the anion redox is triggered. Finally, the series of polyoxyanion $\text{Li}_{(2-x)}\text{TMSiO}_4$ (TM = Mn, Fe, Co, and Ni) compounds

does not exhibit a universal feature of oxygen redox in general but a dependence on the 3d occupation numbers of the TM. This work provides new insights on the anionic redox activity in traditional polyanionic cathode materials and paves the way for future design of high-capacity rechargeable batteries.

■ ASSOCIATED CONTENT

Supporting Information

The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acs.jpcllett.8b02725.

Methods and notes on the reversibility of lithiation/delithiation for the studied compounds, the structural response to the creation of hole polarons during delithiation, sXAS results, the origin of anionic redox activity, the evolution of energy orbitals during delithiation, the average open-circuit voltage during delithiation, and comparison of results between the hybrid functional and Lany–Zunger approach (PDF)

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Notes

The authors declare no competing financial interest.

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